

Optimal non-symmetric Fokker-Planck equation for the convergence to a given equilibrium

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Topic & goals

symmetric Fokker-Planck equation for $f(x, t)$:

- $$f_t = \operatorname{div}(\nabla f + f \nabla V(x)), \quad x \in \mathbb{R}^d, \quad t > 0$$

→ Decay estimate to $f_\infty(x) = c_V e^{-V(x)}$ with rate $\inf_x \lambda_{\min}\left(\frac{\partial^2 V}{\partial x^2}\right)$ by entropy method (Bakry-Emery strategy)

- This rate is sharp for $V(x) = \frac{x^T \mathbf{K}^{-1} x}{2}$, $\mathbf{K} > 0$,
 $f_\infty(x) = c_{\mathbf{K}} e^{-V(x)}$... anisotropic Gaussian

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Theme:

- Non-symmetric perturbations of the drift (that preserve f_∞) can enhance the convergence.
- Goal: Find optimal perturbation that yields best exponential estimate

$$\|f(t) - f_\infty\|_{L^2(f_\infty^{-1})} \leq c e^{-\lambda t} \|f_0 - f_\infty\|_{L^2(f_\infty^{-1})}, \quad t \geq 0,$$

with (1) maximal $\lambda > 0$ and (2) minimal $c \geq 1$.

stochastic applications

- compute expectations w.r.t. measure $\mu_V = e^{-V} dx$ (high dimensions)
- needs to construct an ergodic Markov process with fast convergence to the unique measure μ_V .

References:

- Lelièvre-Nier-Pavliotis: *Optimal non-reversible linear drift for the convergence to equilibrium of a diffusion*, *J. Stat. Phys.*, 2013
- Guillin-Monmarché: *Optimal linear drift for the speed of convergence of an hypoelliptic diffusion*, *Electron. Commun. Probab.*, 2016

Outline:

- 1 Fokker-Planck equations with linear drift: propagator norm
- 2 construction of best (hypocoercive) Fokker-Planck equations
- 3 numerical illustrations
- 4 outlook: Fokker-Planck equations with t -dependent coefficients

degenerate Fokker-Planck equations with linear drift

$$f_t = \operatorname{div} \left(\mathbf{D} \nabla f + \mathbf{C}_x f \right) =: -L f, \quad x \in \mathbb{R}^d \quad (1)$$

with degenerate $0 \leq \mathbf{D} = \mathbf{D}^T \in \mathbb{R}^{d \times d}$ is degenerate parabolic;
(symmetric part of) L is **not coercive**.

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Definition 1 (Villani 2009)

Consider L on Hilbert space H with $\mathcal{K} = \ker L$; let $\tilde{H} \hookrightarrow \mathcal{K}^\perp$ (densely)
(e.g. $H \dots$ weighted L^2 , $\tilde{H} \dots$ weighted H^1).

L is called **hypocoercive** on \tilde{H} if $\exists \lambda > 0, c \geq 1$:

$$\|e^{-Lt} f_0\|_{\tilde{H}} \leq c e^{-\lambda t} \|f_0\|_{\tilde{H}} \quad \forall f_0 \in \tilde{H}$$

- typically $c > 1$

hypo coercive Fokker-Planck equation

$$f_t = \operatorname{div} \left(\mathbf{D} \nabla f + \mathbf{C} x f \right) =: -L f, \quad x \in \mathbb{R}^d \quad (2)$$

Condition A:

- 1 No (nontrivial) subspace of $\ker \mathbf{D}$ is invariant under \mathbf{C}^\top .
(equivalent: L is hypoelliptic.)
 - 2 Let $\mathbf{C} \in \mathbb{R}^{d \times d}$ be positive stable (i.e. $\Re(\lambda^{\mathbf{C}}) > 0$).
[$\Rightarrow \exists$ confinement potential; drift towards $x = 0$.]
- hypoelliptic + confinement = hypo coercive (for FP eq.)

hypo coercive Fokker-Planck equation

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[$\Rightarrow \exists$ confinement potential; drift towards $x = 0$.]
- hypoelliptic + confinement = hypo coercive (for FP eq.)

Lemma 1

Let Condition A hold. Then:

- (2) is hypo coercive.
- $f_\infty(x) = c_{\mathbf{K}} \exp\left(-\frac{x^T \mathbf{K}^{-1} x}{2}\right)$ is the unique steady state, with $\mathbf{K} = \mathbf{K}^T > 0$ solves $2\mathbf{D} = \mathbf{C}\mathbf{K} + \mathbf{K}\mathbf{C}^T$ (continu. Lyapunov eq. for \mathbf{K})

normalization of Fokker-Planck equations

original FP-equation: $f_t = \operatorname{div} \left(\mathbf{D} \nabla_x f + \mathbf{C} x f \right) =: -L f, \quad x \in \mathbb{R}^d$

coordinate transformation: $y := \mathbf{K}^{-1/2} x, \quad g(y) := \sqrt{\det(\mathbf{K})} f(\mathbf{K}^{1/2} y) \Rightarrow$

normalized FP-equation: $g_t = \operatorname{div} \left(\tilde{\mathbf{D}} \nabla_y g + \tilde{\mathbf{C}} y g \right) =: -\tilde{L} g, \quad y \in \mathbb{R}^d$

corresponding drift-ODE: $\frac{d}{dt} y(t) = -\tilde{\mathbf{C}} y(t), \quad \tilde{\mathbf{C}} := \mathbf{K}^{-1/2} \mathbf{C} \mathbf{K}^{1/2}$

with $\tilde{\mathbf{D}} = \mathbf{K}^{-1/2} \mathbf{D} \mathbf{K}^{-1/2} \Rightarrow \tilde{\mathbf{D}} = \tilde{\mathbf{C}}_s \geq 0$

Then $g_\infty(x) = (2\pi)^{-d/2} e^{-|y|^2/2}$.

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L^2 -propagator norm (\rightarrow our main tool):

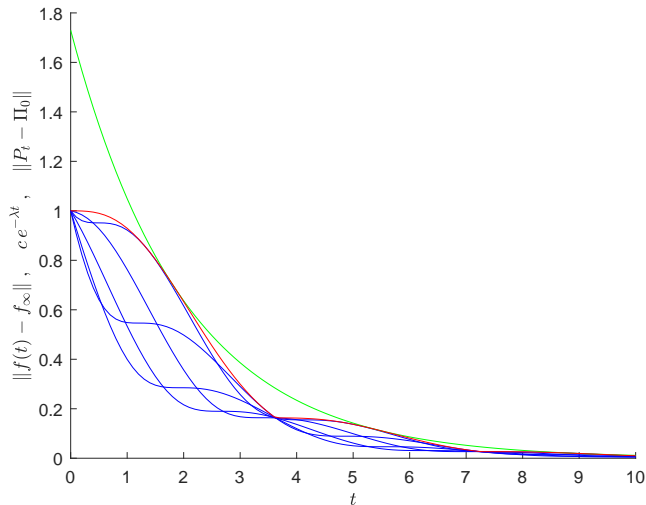
Theorem 2 (AA-Schmeiser-Signorello, CMS 2022)

Let Condition A hold. Then:

$$\left\| e^{-Lt} \right\|_{\mathcal{B}(\{f_\infty\}^\perp)} = \left\| e^{-\tilde{L}t} \right\|_{\mathcal{B}(\{g_\infty\}^\perp)} = \left\| e^{-\tilde{\mathbf{C}}t} \right\|_{\mathcal{B}(\mathbb{R}^d)}, \quad \forall t \geq 0$$

L^2 -propagator norm

propagator norm of Fokker-Planck equ. in 2D: $\|e^{-Lt} - \Pi_0\|_{\mathcal{B}(L^2(\mathbb{R}^2))}$



Proof of Theorem 2; for normalized FP (step 1)

Theorem 3

Let $\tilde{L} = -\operatorname{div}(\tilde{\mathbf{D}} \nabla \cdot + \tilde{\mathbf{C}} y \cdot)$ satisfy Condition A (i.e. \tilde{L} is hypocoercive).

Then

$$\|e^{-\tilde{L}t} - \tilde{\Pi}_0\|_{\mathcal{B}(\mathcal{H})} = \|e^{-\tilde{\mathbf{C}}t}\|_2, \quad t \geq 0.$$

$\tilde{\Pi}_0$... projection on $\operatorname{span}[g_\infty]$, $g_\infty = c e^{-|y|^2/2}$

- \tilde{L} ... non-symmetric. Still, \exists a partially orthogonal decomposition:

$$\mathcal{H} := L^2(g_\infty^{-1}) = \bigoplus_{m \in \mathbb{N}_0}^\perp V^{(m)}; \quad V^{(m)} = \operatorname{span}[g_\alpha(y) := (-1)^{|\alpha|} \nabla^\alpha g_\infty, |\alpha| = m]$$

$$\sigma(L) = \left\{ \sum_{j=1}^d \alpha_j \lambda_j, \alpha \in \mathbb{N}_0^d \right\}; \quad \lambda_j \dots \text{eigenvalues of } \tilde{\mathbf{C}} \in \mathbb{R}^{d \times d}$$

Proof (step 2): evolution in subspaces $V^{(m)}$

$d_\alpha(t)$... coefficient of $g_\alpha(y)$, $\alpha \in \mathbb{N}_0^d$, $y \in \mathbb{R}^d$

ex. $d = 2$:

- $m = 1$: $\frac{d}{dt} \begin{pmatrix} d_{(1,0)} \\ d_{(0,1)} \end{pmatrix} = -\tilde{\mathbf{C}} \begin{pmatrix} d_{(1,0)} \\ d_{(0,1)} \end{pmatrix}$... drift-ODE

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- $m = 2$: $\begin{pmatrix} d_{(2,0)} \\ d_{(1,1)} \\ d_{(0,2)} \end{pmatrix}$... impractical !

better: $D^{(2)}(t) := \begin{pmatrix} d_{(2,0)} & d_{(1,1)}/2 \\ d_{(1,1)}/2 & d_{(0,2)} \end{pmatrix} (t) \in \mathbb{R}^{2 \times 2}$

$$\frac{d}{dt} D^{(2)} = -(\tilde{\mathbf{C}} D^{(2)} + D^{(2)} \tilde{\mathbf{C}}^T)$$

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$$\frac{d}{dt} D^{(2)} = -(\tilde{\mathbf{C}} D^{(2)} + D^{(2)} \tilde{\mathbf{C}}^T)$$

- $m \geq 3$: $D^{(m)}(t)$... symmetric m -order tensor

$$\frac{d}{dt} D^{(m)}(t) = -m \operatorname{Sym} \left(\underbrace{\tilde{\mathbf{C}} \odot D^{(m)}(t)}_{\text{mult. on 1st index}} \right) \quad \dots \quad \text{tensored drift-ODE}$$

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\Rightarrow FP = 2nd quantization of ODE in Bosonic Fock space of \mathbb{R}^2

Proof (step 3): evolution in subspaces $V^{(m)}$

- practical ingredient for estimating the evolution equation in $V^{(m)}$:
rank-1 decomposition of order- m tensors:

$$D^{(m)} = \sum_{k=1}^s \mu_k v_k^{\otimes m}, \quad \mu_k \in \mathbb{R}, v_k \in \mathbb{R}^d \quad (3)$$

minimal s : “symmetric rank” of $D^{(m)}$

Lemma 2

Let (3) be the decomposition of $D^{(m)}(0)$. Then, the evolution in $V^{(m)}$ is given by

$$D^{(m)}(t) = \sum_{k=1}^s \mu_k [v_k(t)]^{\otimes m}, \quad \dot{v}_k = -\tilde{\mathbf{C}} v_k .$$

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- subspaces $V^{(m)}$ are orthogonal
- decay behavior of $\|e^{-\tilde{L}t} - \tilde{\Pi}_0\|_{\mathcal{B}(\mathcal{H})}$ determined only by 1st subspace
→ equivalent to drift-ODE: $\|e^{-\tilde{\mathbf{C}}t}\|_2$

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(non-normalized) Fokker-Planck equations with linear drift

Let steady state $f_{\infty, \mathbf{K}}(x) = \frac{\det(\mathbf{K})^{-1/2}}{(2\pi)^{-d/2}} \exp\left(-\frac{x^T \mathbf{K}^{-1} x}{2}\right)$, $\mathbf{K} > 0$ be given.

Find $f_t = \operatorname{div}(\mathbf{D}\nabla f + \mathbf{C}xf) =: -L_{\mathbf{C}, \mathbf{D}}f$ **with fastest decay**; (4)

diffusion matrix $\mathbf{D} \geq 0$; drift matrix \mathbf{C} : positive stable, i.e. $\Re(\lambda^{\mathbf{C}}) > 0$.

admissible matrices: $\mathcal{I}(\mathbf{K}) := \{(\mathbf{C}, \mathbf{D}) : \mathbf{D} \geq 0, \operatorname{Tr}(\mathbf{D}) \leq d, L_{\mathbf{C}, \mathbf{D}}f_{\infty, \mathbf{K}} = 0\}$

- Without constraint $\operatorname{Tr}(\mathbf{D}) \leq d$ arbitrary decay possible \rightarrow ill-posed.

(non-normalized) Fokker-Planck equations with linear drift

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• Without constraint $\operatorname{Tr}(\mathbf{D}) \leq d$ arbitrary decay possible \rightarrow ill-posed.

Lemma 3 (Guillin-Monmarché 2016)

$$\mathcal{I}(\mathbf{K}) = \{\mathbf{D} \geq 0, \operatorname{Tr}(\mathbf{D}) \leq d; \mathbf{C} = (\mathbf{D} + \mathbf{J})\mathbf{K}^{-1}, \mathbf{J}^T = -\mathbf{J}\}.$$

Lemma 4

Let $\mathbf{K} > 0$, $(\mathbf{C}, \mathbf{D}) \in \mathcal{I}(\mathbf{K})$, \mathbf{C} positive stable.

Then $f_{\infty, \mathbf{K}}$ is the unique (normalized) steady state; (4) is hypocoercive.

Questions

hypocoercivity:
$$\|f(t) - f_{\infty, K}\|_{L^2(f_{\infty, K}^{-1})} \leq c e^{-\lambda t} \|f_0 - f_{\infty, K}\|_{L^2(f_{\infty, K}^{-1})}, \quad t \geq 0 \quad (5)$$

- Q1 Which FP-evolutions converge the fastest, i.e. with **largest rate** λ_{opt} to the steady state in the **operator norm of** $e^{-L_{C,D}t}$ on $\{f_{\infty, K}\}^{\perp} \subset \mathcal{H} := L^2(\mathbb{R}^d, f_{\infty, K}^{-1})$?
- Q2 When the best decay rate is fixed, what is the **infimum of the multiplicative constant**, c_{inf} , in the decay estimate (5)?
- Q3 For a fixed $K > 0$ and the corresponding λ_{opt} , and for any $c > c_{inf}$, **which pair(s) of matrices** $(\mathbf{C}_{opt}(c), \mathbf{D}_{opt}(c) \geq 0)$ yields the convergence estimate (5) with the constants (λ_{opt}, c) ?
- Q4 For such an optimal pair of matrices, what **bound on** $\|\mathbf{C}_{opt}\|$ can be found, and how does this bound grow w.r.t. to the space dimension d ?

Results from the literature

Lemma 5 (Guillin-Monmarché 2016)

- *Question Q1: $\lambda_{opt} = \lambda_{max}(\mathbf{K}^{-1})$.*
- *Question Q2: They can only reach multiplicative constants $c > \sqrt{\kappa(\mathbf{K})} e$.*
- *Questions Q3+Q4: Their drift matrix grows like $\|\mathbf{C}_{opt}\| = \mathcal{O}(d^2)$ (with piecewise constant coefficients).*

$\kappa(\mathbf{K})$... condition number

Strategy for improvement

non-symmetric FP-equation:

$$f_t = \operatorname{div}(\mathbf{D}\nabla f + \mathbf{C}xf) =: -L_{\mathbf{C},\mathbf{D}}f$$

corresponding drift-ODE:

$$\frac{d}{dt}y(t) = -\tilde{\mathbf{C}}y(t), \quad \tilde{\mathbf{C}} := \mathbf{K}^{-1/2}\mathbf{C}\mathbf{K}^{1/2}$$

Main tool:

Theorem 4 (AA-Schmeiser-Signorello, CMS 2022)

Let $\mathbf{K} > 0$, $(\mathbf{C}, \mathbf{D}) \in \mathcal{I}(\mathbf{K})$, \mathbf{C} positive stable. Then:

$$\left\| e^{-L_{\mathbf{C},\mathbf{D}}t} \right\|_{\mathcal{B}(\{f_{\infty,\mathbf{K}}\}^{\perp})} = \left\| e^{-\tilde{\mathbf{C}}t} \right\|_{\mathcal{B}(\mathbb{R}^d)}, \quad \forall t \geq 0,$$

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- This reduces the PDE-optimization problem to an ODE-problem, and allows for sharp result.
- Replacing a hypocoercive entropy method used in [Guillin-Monmarché 2016]; block-diagonal decomposition of the FP-propagator used in [Lelièvre-Nier-Pavliotis 2013].

Main result: optimal constants

Theorem 5 (AA-Signorello 2021)

Let $\mathbf{K} > 0$ be given. Then:

- (a) Questions Q2+Q3: $c_{inf} = 1$. For any constant $c > 1$, there exists a pair $(\mathbf{C}_{opt}(c), \mathbf{D}_{opt}(c)) \in \mathcal{I}(\mathbf{K})$ such that

$$\left\| e^{-L_{\mathbf{C}_{opt}, \mathbf{D}_{opt}} t} \right\|_{\mathcal{B}(V_0^\perp)} \leq c e^{-\max(\sigma(\mathbf{K}^{-1}))t}, \quad t \geq 0. \quad (6)$$

- (b) Question Q4: The matrices from (a) satisfy

$$\|\mathbf{C}_{opt}\|_{\mathcal{F}} \leq \lambda_{opt} \left[d + \sqrt{\kappa(\mathbf{K})} \frac{2\pi c^2}{\sqrt{3}(c^2 - 1)} \sqrt{d} (d-1) \right], \quad \|\mathbf{D}_{opt}\|_{\mathcal{F}} = d. \quad (7)$$

$\|\mathbf{C}\|_{\mathcal{F}}$... Frobenius norm

Proof-idea (refinement of [Lelièvre-Nier-Pavliotis 2013], [Guillin-Monmarché 2016])

① normalize FP-equation: $y := \mathbf{K}^{-1/2}x$

$$\Rightarrow g_t = \operatorname{div}(\tilde{\mathbf{D}}\nabla g + \tilde{\mathbf{C}}yg), \quad g_\infty(y) = (2\pi)^{-d/2}e^{-|y|^2/2}, \quad y \in \mathbb{R}^d,$$

with $\tilde{\mathbf{D}} := \mathbf{K}^{-1/2}\mathbf{D}\mathbf{K}^{-1/2} \geq 0$, $\tilde{\mathbf{J}} := \mathbf{K}^{-1/2}\mathbf{J}\mathbf{K}^{-1/2} = -\tilde{\mathbf{J}}^T$,
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② Construction of \mathbf{D} :

Maximal decay rate of $e^{-\tilde{\mathbf{C}}t}$, $\lambda_{opt} = \lambda_{max}(\mathbf{K}^{-1})$ only possible if
 $\operatorname{ran}(\mathbf{D}) \subset \operatorname{eigenspace}_{\lambda_{max}}(\mathbf{K}^{-1})$;
e.g. $\mathbf{D} := d v \otimes v$ (with $\mathbf{K}^{-1}v = \lambda_{opt}v$, $\|v\| = 1$) ... rank 1.

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③ \exists algorithmic construction of $\tilde{\mathbf{J}}, \mathbf{P} > 0$ such that:

$$\mathbf{P} \underbrace{(\tilde{\mathbf{D}} + \tilde{\mathbf{J}})}_{=\tilde{\mathbf{C}}} + \underbrace{(\tilde{\mathbf{D}} - \tilde{\mathbf{J}})}_{=\tilde{\mathbf{C}}^T} \mathbf{P} = 2\lambda_{opt}\mathbf{P} \dots \text{continuous Lyapunov equation for } \mathbf{P}$$

Proof-idea (cont'd)

④ decay of drift-ODE $\dot{y}(t) = -\tilde{\mathbf{C}}y(t)$ in norm $\|y\|_{\mathbf{P}}^2 := \langle y, \mathbf{P}y \rangle$:

$$\frac{d}{dt} \|y(t)\|_{\mathbf{P}}^2 = -\langle y(t), [\mathbf{P}\tilde{\mathbf{C}} + \tilde{\mathbf{C}}^T \mathbf{P}]y(t) \rangle = -2\lambda_{opt} \|y(t)\|_{\mathbf{P}}^2.$$

$$\Rightarrow \|y(t)\|_{\mathbf{P}} = e^{-\lambda_{opt}t} \|y(0)\|_{\mathbf{P}}, \quad t \geq 0.$$

Rem: $\tilde{\mathbf{C}}$ is not coercive $\rightarrow \mathbf{P}$ provides the “hypocoercivity norm.”

In Euclidean matrix norm:

$$\|e^{-\tilde{\mathbf{C}}_{opt}t}\|_{\mathcal{B}(\mathbb{R}^d)} \leq \sqrt{\kappa(\mathbf{P})} e^{-\lambda_{opt}t}, \quad t \geq 0,$$

$\kappa(\mathbf{P})$... condition number; can be chosen arbitrarily close to 1 with a “good” choice of $\tilde{\mathbf{C}}$. □

Proof-idea (cont'd)

④ decay of drift-ODE $\dot{y}(t) = -\tilde{\mathbf{C}}y(t)$ in norm $\|y\|_{\mathbf{P}}^2 := \langle y, \mathbf{P}y \rangle$:

$$\frac{d}{dt} \|y(t)\|_{\mathbf{P}}^2 = -\langle y(t), [\mathbf{P}\tilde{\mathbf{C}} + \tilde{\mathbf{C}}^T \mathbf{P}]y(t) \rangle = -2\lambda_{opt} \|y(t)\|_{\mathbf{P}}^2.$$

$$\Rightarrow \|y(t)\|_{\mathbf{P}} = e^{-\lambda_{opt} t} \|y(0)\|_{\mathbf{P}}, \quad t \geq 0.$$

Rem: $\tilde{\mathbf{C}}$ is not coercive $\rightarrow \mathbf{P}$ provides the “hypo-coercivity norm.”

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Remarks:

- $(\mathbf{C}_{opt}(c), \mathbf{D}_{opt}(c))$ is not unique; $\tilde{\mathbf{C}}_{opt}^T(c)$ yields an alternative.
- This optimal FP eq. has maximal *hypo-coercivity index* $d - 1$.

Outline:

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- 2 construction of best (hypocoercive) Fokker-Planck equations
- 3 numerical illustrations
- 4 outlook: Fokker-Planck equations with t -dependent coefficients

Numerical illustration (1)

2D Example:

- Given covariance matrix $\mathbf{K} = \text{diag}(\frac{1}{\varepsilon}, 1)$, $\varepsilon = 0.05$.
 $\Rightarrow \lambda_{opt} = \lambda_{max}(\mathbf{K}^{-1}) = 1$.
- For any $c > 1$ in the decay estimate $\|f(t) - f_{\infty, K}\|_{L^2(f_{\infty, K}^{-1})} \leq c e^{-\lambda_{max} t}$:

$$\Rightarrow \mathbf{D}_{opt} = \tilde{\mathbf{D}}_{opt} = \begin{pmatrix} 0 & 0 \\ 0 & 2 \end{pmatrix}$$

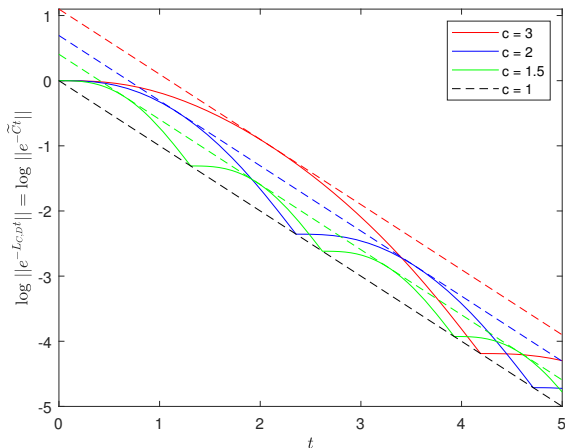
$$\mathbf{C}_{opt}(c) = \begin{pmatrix} 0 & -\frac{\mu}{\sqrt{\varepsilon}} \\ \sqrt{\varepsilon}\mu & 2 \end{pmatrix}, \quad \tilde{\mathbf{C}}_{opt}(c) = \begin{pmatrix} 0 & -\mu \\ \mu & 2 \end{pmatrix}, \quad \mu := \frac{c^2 + 1}{c^2 - 1}.$$

- $c \searrow 1 \Rightarrow \mu \rightarrow \infty \dots$ high-rotational limit

Practical tradeoff:


better convergence vs. smaller matrix \mathbf{C} (\rightarrow allows for larger time steps in the numerics of the mentioned Markov process)

Numerical illustration of $\|f(t) - f_{\infty,K}\|_{L^2(f_{\infty,K}^{-1})} \leq c e^{-\lambda_{\max} t}$



— exact propagator norms of FP-equation and its drift-ODE (for $c = 3$):

$$\left\| e^{-L_C, D t} \right\|_{\mathcal{B}(\{f_{\infty,K}\}^\perp)} = \left\| e^{-\tilde{C}t} \right\|_{\mathcal{B}(\mathbb{R}^d)}, \quad t \geq 0$$

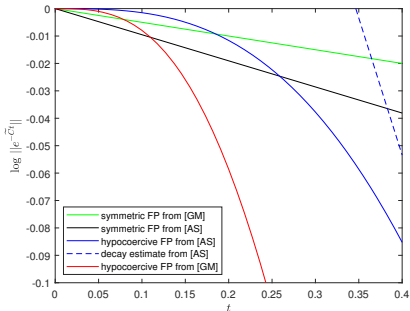
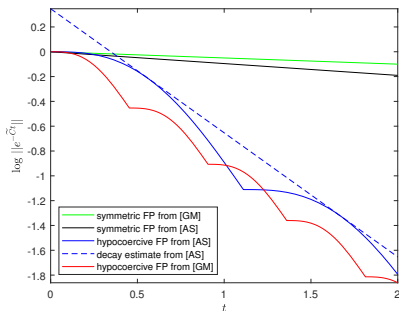
- - - corresponding exponential decay estimate (as optimal upper envelop) 

Numerical illustration: comparison to [Guillin-Monmarché]

- Given covariance matrix $\mathbf{K} = \text{diag}(\frac{1}{\varepsilon}, 1)$, $\varepsilon = 0.05$, given $c = \sqrt{2}$.

$$\Rightarrow \tilde{\mathbf{D}}_{opt} = \begin{pmatrix} 0 & 0 \\ 0 & 2 \end{pmatrix}, \quad \tilde{\mathbf{c}}_{opt}^{AS} = \begin{pmatrix} 0 & -3 \\ 3 & 2 \end{pmatrix}, \quad \tilde{\mathbf{c}}_{opt}^{GM} = \begin{pmatrix} 0 & -7 \\ 7 & 2 \end{pmatrix}$$

- Estimate in [GM] is not sharp \rightarrow more rotation than “necessary” for the bound $c e^{-\lambda_{max} t}$ (---) is used \rightarrow unfavorable for time step restriction.

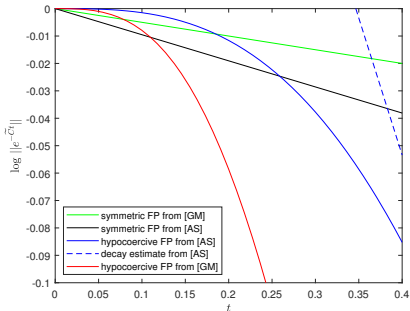
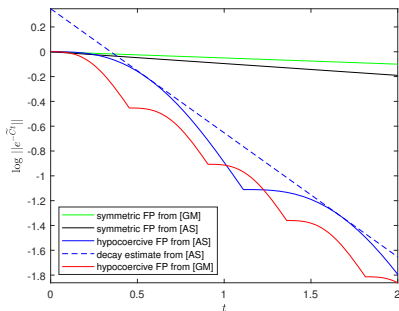


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- zoom (at $t \approx 0$) on right: symmetric FP-evolution decays initially faster!

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FP with t -dependent coefficients; \mathbf{K} ... t -independent

Let steady state $f_{\infty, \mathbf{K}}(x) = \frac{\det(\mathbf{K})^{-1/2}}{(2\pi)^{-d/2}} \exp\left(-\frac{x^T \mathbf{K}^{-1} x}{2}\right)$, $\mathbf{K} > 0$ be given.

$$f_t = \operatorname{div}(\mathbf{D}(t)\nabla f + \mathbf{C}(t)xf), \quad (\mathbf{C}(t), \mathbf{D}(t)) \in \mathcal{I}(\mathbf{K}) \quad \forall t \geq 0.$$

corresponding drift-ODE: $\frac{d}{dt}y(t) = -\tilde{\mathbf{C}}(t)y(t)$, $\tilde{\mathbf{C}}(t) := \mathbf{K}^{-\frac{1}{2}}\mathbf{C}(t)\mathbf{K}^{\frac{1}{2}}$

- A split FP-evolution yielded in [Guillin-Monmarché] a significant improvement of the decay estimate, and enabled $\|\mathbf{C}_{opt}\|_{\mathcal{F}} = \mathcal{O}(d^2)$:

$$\begin{cases} f_t = \operatorname{div}(\nabla f + \mathbf{K}^{-1}xf), & 0 \leq t \leq t_0, & \text{symmetric FP,} \\ f_t = \operatorname{div}(\mathbf{D}_{opt}\nabla f + \mathbf{C}_{opt}xf), & t > t_0, & \text{non-symm. FP.} \end{cases}$$

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- Q5: Does the symmetric FP-evolution on $[0, t_0]$ give a *true* improvement for the FP-propagator norm or 'only' a better analytic estimate?

This was difficult to decide for the PDE so far.

FP with t -dependent coefficients

Main tool:

Theorem 6 (AA-Signorello 2021)

Let $\mathbf{K} > 0$, $(\mathbf{C}(t), \mathbf{D}(t)) \in \mathcal{I}(\mathbf{K})$, $\mathbf{C}(t)$ positive stable $\forall t \geq 0$. Then:

$$\|S(t_2, t_1)\|_{\mathcal{B}(\{f_{\infty, \mathbf{K}}\}^\perp)} = \|T(t_2, t_1)\|_{\mathcal{B}(\mathbb{R}^d)}, \quad \forall 0 \leq t_1 \leq t_2 < \infty$$

- FP-propagator $S(t_2, t_1)$ maps $f(t_1) \in L^2(f_{\infty, \mathbf{K}}^{-1})$ to $f(t_2)$.
- propagator of drift-ODE $\dot{y} = -\tilde{\mathbf{C}}(t)y$: $T(t_2, t_1)$ maps $y(t_1) \in \mathbb{R}^d$ to $y(t_2)$.

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Proof-idea:

- All FP-equations with frozen t have the same steady state $f_{\infty, K}(x)$.
- \Rightarrow The FP-normalizations, the space $L^2(f_{\infty, K}^{-1})$, and the subspace decomposition $\mathcal{H} := L^2(g_\infty^{-1}) = \bigoplus_{m \in \mathbb{N}_0}^\perp V^{(m)}$ are all t -independent.

FP with t -dependent coefficients: 2D numerical case study

Remark on Q5: $\|e^{-\tilde{\mathbf{C}}t}\|_{\mathcal{B}(\mathbb{R}^d)} = 1 - \lambda_{\min}(\tilde{\mathbf{C}}_s) t + \mathcal{O}(t^2)$ as $t \rightarrow 0$

\Rightarrow An initially symmetric FP-evolution always decays faster than a hypocoercive FP-evolution ($\tilde{\mathbf{C}}_s = \tilde{\mathbf{D}}$; typically $\text{rank}(\tilde{\mathbf{D}}) = 1$).

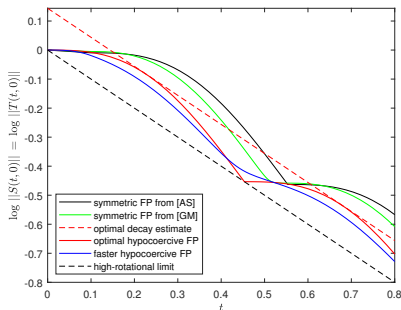
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But it backfires later! Even when switching to **opt. non-symmetric FP later**.

- Given covariance matrix $\mathbf{K} = \text{diag}(\frac{1}{\varepsilon}, 1)$, $\varepsilon = 0.05$, given $c := \sqrt{4/3}$.



ref. case: $\tilde{\mathbf{C}}_1 = [0 \ -7; 7 \ 2] \ \forall t \geq 0$.

split FP-evolutions:

symm. FP on $t \in [0, 0.1]$, then $\tilde{\mathbf{C}}_1$.

faster rot. with $\tilde{\mathbf{C}}_2 = [0 \ -11; 11 \ 2]$
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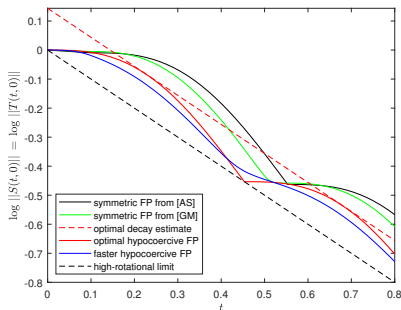
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- Open question: What is the best $\mathbf{C}(t)$, $\mathbf{D}(t)$?

Conclusion

- Construction of Fokker-Planck equations $f_t = \text{div}(\mathbf{D}\nabla f + \mathbf{C}xf)$ with optimal decay.
- main tool: Propagator norms of FP-equation and corresponding drift-ODE ($\dot{x} = -\tilde{\mathbf{C}}x$) coincide.
- t -dependent coefficients can enhance the decay of Fokker-Planck equations.

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References

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Thanks for
your attention